

BOUNDS FOR THE LARGEST EIGENVALUE IN ALTERNANT HYDROCARBONS
A TOPOLOGICAL APPROACH

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Abstract. Certain topological upper and lower bounds for the largest eigenvalue of alternant hydrocarbons are derived.

1. Introduction

The eigenvalue-eigenvector problem of the adjacency matrix of a graph has attracted the attention of graph theoreticians over a considerable period of time¹. For certain classes of graphs, within the framework of Hückel Molecular Orbital (HMO) theory, the graph spectrum can be solved in a closed analytical form, most frequently using the symmetry properties of a graph². Thus, graph spectral-structural analysis was carried out analytically for rings³, chains³, stars and other trees^{4,5,6} etc.

The interval in which the eigenvalues X_j of a graph lie is limited and given by Frobenius theorem^{2,7},

$$-d_{\max} \leq X_j \leq d_{\max} \quad j = 1, 2, \dots, N$$

where, d_{\max} is the maximal vertex degree in a graph. Therefore, for the Hückel graphs^{1,2} the whole spectrum lies in the interval from -3 to +3. In the linear polyenes and annulenes the eigenvalues lie in the interval from -2 to +2. The largest and the smallest eigenvalues of a molecular graph have been already investigated in some detail.^{8,9} In the present study, topologically dependent alternatives of the abovementioned bounds are given.

2. Theory.

Suppose, G is an undirected planar graph, having r rings, e edges and N vertices such that the degree (d_i) of every vertex in G is $1 \leq d_i \leq 3$ (i.e. G is a Hückel graph^{1,2}). Furthermore, let N be an even number and $P(X)$ be the corresponding characteristic polynomial,

$$P(X) = X^N + a_1 X^{N-1} + a_2 X^{N-2} + \dots + a_{N-1} X + a_N \quad (1)$$

For alternant hydrocarbons the graph eigenvalues appear in pairs $\bar{X}_1, \bar{X}_2, \dots, \bar{X}_{N/2}$ and $P(X)$ can be written in the form of,

$$P(X) = (X^2 - X_1^2)(X^2 - X_2^2) \dots (X^2 - X_{N/2}^2) \quad (2)$$

We assume : $X_1 \geq X_2 \geq \dots \geq X_{N/2}$. The quantity we are concerned with in the present paper is the largest eigenvalue, X_1 . Within the HMO theory X_1 corresponds to the

energy of the lowest occupied molecular orbital. It is easy to see that the following equations hold^{10,11,12},

$$e = \sum_{i=1}^{N/2} x_i^2 \quad (3)$$

$$a_4 = \sum_{i < j}^{N/2} x_i^2 x_j^2 \quad (4)$$

$$e^2 = \sum_{i=1}^{N/2} x_i^4 + 2a_4 \quad (5)$$

$$a_6 = - \sum_{i < j < k}^{N/2} x_i^2 x_j^2 x_k^2 \quad (6)$$

$$a_8 = \sum_{i < j < k < l}^{N/2} x_i^2 x_j^2 x_k^2 x_l^2 \quad (7)$$

On the other hand, the following holds ,

$$\sum_{i \neq j \neq k}^{N/2} x_i^4 x_j^2 x_k^2 = \sum_{i=1}^{N/2} x_i^4 (a_4 - x_i^2 (e - x_i^2)) \quad (8)$$

$$\sum_{i \neq j \neq k}^{N/2} x_i^4 x_j^2 x_k^2 = a_4 \sum_{i=1}^{N/2} x_i^4 - e \sum_{i=1}^{N/2} x_i^6 + \sum_{i=1}^{N/2} x_i^8 \quad (9)$$

By using eq.5 , eq.9 is changed into,

$$\sum_{i \neq j \neq k}^{N/2} x_i^4 x_j^2 x_k^2 = a_4 (e^2 - 2a_4) - e \sum_{i=1}^{N/2} x_i^6 + \sum_{i=1}^{N/2} x_i^8 \quad (10)$$

Now, squaring both sides of eq.4 one gets,

$$a_4^2 = \sum_{i < j}^{N/2} x_i^4 x_j^4 + \frac{1}{2} \sum_{i \neq j \neq k}^{N/2} x_i^4 x_j^2 x_k^2 + \frac{1}{6} \sum_{i \neq j \neq k \neq l}^{N/2} x_i^2 x_j^2 x_k^2 x_l^2 \quad (11)$$

Inserting eqs.7 and 10 into eq.11 and noting that¹¹,

$$\sum_{i=1}^{N/2} x_i^6 = e^3 - 3ea_4 + 3a_6 \quad (12)$$

one obtains,

$$\sum_{i < j}^{N/2} x_i^4 x_j^4 = 5a_4^2 - 8ea_4 + 2e^4 + 6ea_6 - 2 \sum_{i=1}^{N/2} x_i^8 - 6a_8 \quad (13)$$

Multiplying eqs.3 and 6 side by side, one gets,

$$ea_6 = \sum_{i=1}^{N/2} x_i^4 (a_4 - x_i^2(e - x_i^2)) + 4a_8 \quad (14)$$

$$ea_6 = a_4 \sum_{i=1}^{N/2} x_i^4 - e \sum_{i=1}^{N/2} x_i^6 + \sum_{i=1}^{N/2} x_i^8 + 4a_8 \quad (15)$$

Using eqs.5 ,12 ,15 and one has,

$$\sum_{i=1}^{N/2} x_i^8 = 4ea_6 - 4e^2 a_4 + 2a_4^2 - 4a_8 + e^4 \quad (16)$$

Then eq.13 becomes,

$$\sum_{i < j}^{N/2} x_i^4 x_j^4 = a_4^2 - 2ea_6 + 2a_8 \quad (17)$$

On the other hand, it is evident that, $x_1^8 < \sum_{i=1}^{N/2} x_i^8$ and $x_1^4 <$

$\sum_{i=1}^{N/2} x_i^4$. Then by using eqs.16 and 5 one gets,

$$x_1 < \sqrt[8]{a^4 + 4ea_6 + 2a_4^2 - 4e^2a_4 - 4a_8} \equiv x_U \quad (18)$$

$$x_1 < \sqrt[4]{e^2 - 2a_4} \quad (19)$$

The right hand side of inequalities 18 and 19 represent two different upper bounds of x_1 of which the former one is denoted by x_U .

Now, consider eq.17. Since, $\sum_{i < j}^{N/2} x_i^4 x_j^4 > \sum_{i=2}^{N/2} x_1^4 x_i^4$,

$$\sum_{i=2}^{N/2} x_1^4 x_i^4 < a_4^2 - 2ea_6 + 2a_8 \quad (20)$$

Note that, $\sum_{i=2}^{N/2} x_i^4 = \sum_{i=1}^{N/2} x_i^4 - x_1^4$. Then ineq.20 becomes,

$$x_1^4 \left(\sum_{i=1}^{N/2} x_i^4 - x_1^4 \right) < a_4^2 - 2ea_6 + 2a_8 \quad (21)$$

and because of eq.5 one obtains

$$(e^2 - 2a_4) x_1^4 - x_1^8 < a_4^2 - 2ea_6 + 2a_8 \quad (22)$$

Ineq.22 can be rearranged as,

$$x_1^8 - (e^2 - 2a_4) x_1^4 + (a_4^2 - 2ea_6 + 2a_8) \geq 0 \quad (23)$$

what implies ,

$$X_1 > \sqrt[4]{0.5(e^2 - 2a_4) + \sqrt{(e^4 + 8ea_6 - 8a_8 - 4e^2a_4)}} \quad (24)$$











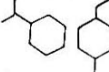



Obviously, the right hand side of the above inequality represents a lower bound, X_L for X_1 . Note that, X_L is an upper bound for X_2 .

3. Conclusion.

The upper and lower bounds for eigenvalues may be found by using various topologically irrelevant mathematical approaches¹³. The topological bounds X_U and X_L derived in the present work estimate the upper and lower bounds of the largest eigenvalue X_1 of alternant hydrocarbons much better (Table 1). The other upper bound expressed by ineq.19 is much easier to evaluate than X_U (ineq.18) but it is obviously less sensitive to topological changes occurring in the molecular structures. Since, the coefficients a_6 and a_8 may reflect the fine details of the topology of a given system, X_U and X_L values should generally be different for isomeric alternant hydrocarbons (see Table 1).

TABLE 1. The Bounds For The Largest Eigenvalue, X_1 , of Various Alternant Systems.

The lower bound, X_L , is given by eq.18 and the upper bound, X_U by eq.24. The X_1 values are excerpted from Ref.14.

Molecule	a_2	a_4	a_6	a_8	X_1	X_U	X_L
	-3	1	0	0	1.618	1.618	1.618
	-5	6	-1	0	1.802	1.813	1.801
	-7	13	-7	1	2.095	2.103	2.093
	-9	26	-30	13	2.170	2.183	2.151
	-6	9	-4	0	2.000	2.001	1.997
	-7	13	-7	0	2.101	2.106	2.096
	-8	19	-16	4	2.136	2.146	2.127
	-8	18	-13	1	2.194	2.199	2.185
	-8	18	-12	0	2.175	2.185	2.165
	-8	18	-12	1	2.170	2.183	2.162
	-9	25	-25	8	2.193	2.214	2.178
	-10	33	-44	24	2.214	2.243	2.181
	-9	22	-16	1	2.355	2.365	2.350
	-11	41	-65	43	2.303	2.322	2.259

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