

## On Almost–Equienergetic Graphs

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(Received March 17, 2013)

### Abstract

Two graphs are said to be equienergetic if their energies are equal. In the paper *MATCH Commun. Math. Comput. Chem.* **61** (2009) 451–461 the concept of *almost-equienergetic* graphs was put forward, based on the observation that in some cases the (non-zero) difference between the energies of two graphs is very small. We now estimate the minimal value of this difference.

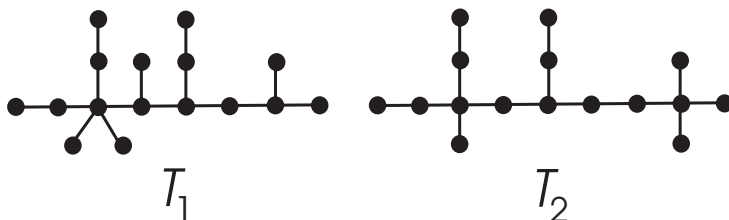
## 1 Introduction

Let  $G$  be a graph of order  $n$  and let its eigenvalues (i.e., the eigenvalues of the  $(0, 1)$ -adjacency matrix of  $G$ ) be  $\lambda_1, \lambda_2, \dots, \lambda_n$ . The *energy* of  $G$  is defined as

$$E(G) = \sum_{i=1}^n |\lambda_i| .$$

For details on the theory of graph energy see the reviews [7, 8, 10], the book [17], and the references cited therein. Two graphs  $G_a$  and  $G_b$  are said to be *equienergetic* if the condition  $E(G_a) - E(G_b) = 0$  is satisfied. This concept was introduced in 2004, independently by Balakrishnan [1] and Brankov et al. [2]. Since then, numerous pairs, triplets, and larger families of equienergetic graphs have been discovered and/or constructed [5, 9, 11–16, 18–23].

Performing a computer-aided search for equienergetic trees [18], it was noticed that there exist pairs of trees for which the difference  $E(G_a) - E(G_b)$  is remarkably small. A characteristic example of this kind is depicted in Fig. 1.



**Fig. 1.** Two trees whose energies differ only slightly:  $E(T_1) = 18.090756640280765\dots$ ,  $E(T_2) = 18.090756641775140\dots$

Based on this observation, the concept of *almost-equienergetic* graphs was conceived [18]. However, a rigorous definition of almost-equienergeticity could not be given. In [18] we read:

*... There also exist trees whose energies are different, but remarkably close. These we refer to as almost-equienergetic. ... What “remarkably close” means for the energy of two graphs is a theme for debate. ... We tentatively and to a great degree arbitrarily call two graphs  $G_a$  and  $G_b$  almost-equienergetic if  $0 < |E(G_a) - E(G_b)| < 10^{-8}$ .*

In this note we show that the limit value  $10^{-8}$  is indeed arbitrary and unjustified, and offer arguments in favor of the possibility that the difference  $E(G_a) - E(G_b)$  can become much smaller.

## 2 Preparatory considerations

Throughout this paper we will consider bipartite graphs. The characteristic polynomial of a bipartite graph  $G$  of order  $n$  is of the form [3]

$$\phi(G, \lambda) = \sum_{k=0}^{\lfloor n/2 \rfloor} (-1)^k b(G, k) \lambda^{n-2k} \tag{1}$$

where  $b(G, 0) = 1$  and  $b(G, k) \geq 0$  for all  $k$ ,  $1 \leq k \leq \lfloor n/2 \rfloor$ .

Let thus  $G_a$  and  $G_b$  be two bipartite graphs of order  $n_a$  and  $n_b$ , respectively. According to a classical result by Coulson and Jacobs [4], if  $n_a = n_b$ , then

$$E(G_a) - E(G_b) = \frac{2}{\pi} \int_0^{+\infty} \ln \frac{\phi(G_a, ix)}{\phi(G_b, ix)} dx$$

where  $i = \sqrt{-1}$ . If  $n_a < n_b$ , then a slightly modified form of the above integral expression applies:

$$E(G_a) - E(G_b) = \frac{2}{\pi} \int_0^{+\infty} \ln \frac{(ix)^{n_b-n_a} \phi(G_a, ix)}{\phi(G_b, ix)} dx .$$

In both cases, in view of Eq. (1),

$$E(G_a) - E(G_b) = \frac{2}{\pi} \int_0^{+\infty} \ln \frac{\sum_{k \geq 0} b(G_a, k) x^{2(m-k)}}{\sum_{k \geq 0} b(G_b, k) x^{2(m-k)}} dx \tag{2}$$

where  $m = \lfloor \frac{1}{2} \max\{n_a, n_b\} \rfloor$ .

Bearing in mind Eq. (2), we see that without loss of generality it may be assumed that the graphs  $G_a$  and  $G_b$  have equal number  $n$  of vertices, and that  $n = 2m$ . If so, then replacing  $b(G_a, k)$  by  $a_{2k}$  and  $b(G_b, k)$  by  $b_{2k}$ , we get

$$E(G_a) - E(G_b) = \frac{2}{\pi} \int_0^{+\infty} \ln \frac{P(x)}{Q(x)} dx \tag{3}$$

with

$$P(x) = x^n + a_2 x^{n-2} + a_4 x^{n-4} + \dots + a_n = \sum_{k=0}^m a_{2k} x^{n-2k} , \quad a_0 = 1$$

$$Q(x) = x^n + b_2 x^{n-2} + b_4 x^{n-4} + \dots + b_n = \sum_{k=0}^m b_{2k} x^{n-2k} , \quad b_0 = 1 .$$

In what follows, by investigating the integral

$$\mathcal{I} = \int_0^{+\infty} \ln \frac{P(x)}{Q(x)} dx \tag{4}$$

we establish some results relevant for the almost-equienergeticity concept.

### 3 Main result

Let  $n = 2m$ ,  $m \in \mathbb{N}$ , and

$$P(x) = x^n + a_2 x^{n-2} + a_4 x^{n-4} + \dots + a_n = \sum_{k=0}^m a_{2k} x^{n-2k}, \quad a_0 = 1$$

$$Q(x) = x^n + b_2 x^{n-2} + b_4 x^{n-4} + \dots + b_n = \sum_{k=0}^m b_{2k} x^{n-2k}, \quad b_0 = 1$$

where  $a_i, b_i \in \mathbb{N}_0 = \mathbb{N} \cup \{0\}$  and  $b_i \leq a_i$ . Let  $\mathcal{I}$  be given by Eq. (4).

Since  $\lim_{n \rightarrow +\infty} \frac{P(x)}{Q(x)} = 1$  and  $\ln x$  is a continuous function on  $(0, +\infty)$ , integration by parts yields

$$\begin{aligned} \mathcal{I} &= x \ln \frac{P(x)}{Q(x)} \Big|_0^{+\infty} - \int_0^{+\infty} \frac{P'(x)Q(x) - Q'(x)P(x)}{P(x)Q(x)} x dx \\ &= \int_0^{+\infty} \frac{P(x)Q'(x) - P'(x)Q(x)}{P(x)Q(x)} x dx. \end{aligned} \tag{5}$$

Choose now the polynomials  $P$  and  $Q$  so that  $a_i = b_i$  for all  $i \neq n - i_0$  and  $a_{n-i_0} = b_{n-i_0} + \ell$ , for some  $\ell \in \mathbb{N}$ . Then

$$P(x) = Q(x) + \ell x^{i_0} \quad \text{and} \quad P'(x) = Q'(x) + \ell i_0 x^{i_0-1}$$

and therefore

$$P(x)Q'(x) - P'(x)Q(x) = \ell x^{i_0-1} (xQ'(x) - i_0 Q(x)).$$

Substituting this back into (5), we get

$$\mathcal{I}_1 = \ell \int_0^{+\infty} \frac{x^{i_0} (xQ'(x) - i_0 Q(x))}{Q(x)(Q(x) + \ell x^{i_0})} dx.$$

Consider now the polynomials  $\tilde{P}(x)$  and  $\tilde{Q}(x)$ , such that  $\tilde{P}(x) = P(x) + \ell x^{i_0}$  and  $\tilde{Q}(x) = Q(x) + \ell x^{i_0}$ , i.e.,  $\tilde{P}(x) = Q(x) + 2\ell x^{i_0}$  and  $\tilde{Q}(x) = Q(x) + \ell x^{i_0}$ . Then the corresponding integral is

$$\mathcal{I}_2 = \ell \int_0^{+\infty} \frac{x^{i_0} (xQ'(x) - i_0 Q(x))}{(Q(x) + x^{i_0})(Q(x) + 2\ell x^{i_0})} dx.$$

Evidently,  $\mathcal{I}_2 < \mathcal{I}_1$ .

Continuing this procedure, namely by increasing the coefficients  $a_{i_0}$  and  $b_{i_0}$  each time by  $\ell \geq 1$ , we get a decreasing sequence of integrals:

$$\mathcal{I}_t = \ell \int_0^{+\infty} \frac{x^{i_0} (x Q'(x) - i_0 Q(x))}{(Q(x) + (t-1)x^{i_0})(Q(x) + t\ell x^{i_0})} dx \quad , \quad t = 1, 2, 3, \dots .$$

In what follows we demonstrate that  $\mathcal{I}_t$  tends to zero as  $t \rightarrow \infty$ .

Let  $n$ , the degree of the polynomials be fixed, and let the coefficients at  $x^{i_0}$  differ by  $\ell \in \mathbb{N}$ , i.e.,  $b_{n-i_0} + \ell = a_{n-i_0}$ . For the sake of simplicity, we denote  $b_{n-i_0} = b$ . We show that for any  $\varepsilon > 0$ , the coefficient  $b$  can be determined so that the value of the integral (4) be less than  $\varepsilon$ .

**Lemma 1.** For  $k \geq 2$ ,

$$\int_0^{+\infty} \frac{dx}{b+x^k} = b^{1/k-1} \frac{\pi}{k} \operatorname{csc} \frac{\pi}{k} . \tag{6}$$

*Proof.* For  $\Re \nu > \Re \mu > 0$  it is known [6, formula 3.241. 2,p. 319] that

$$\int_0^{+\infty} \frac{x^{\mu-1} dx}{1+x^\nu} = \frac{\pi}{\nu} \operatorname{csc} \frac{\mu\pi}{\nu} . \tag{7}$$

In our case,  $k \geq 2 > \mu = 1 > 0$ , and so the integral on the left-hand side of (6) is transformed as:

$$\int_0^{+\infty} \frac{dx}{b+x^k} = \frac{1}{b} \int_0^{+\infty} \frac{dx}{1+(b^{-1/k}x)^k} = \frac{b^{1/k}}{b} \int_0^{+\infty} \frac{dt}{1+t^k} .$$

The right-hand side of (6) follows now directly from (7). □

Recall that  $\operatorname{csc} x = 1/\sin x$ .

**Theorem 1.** Let  $\ell \in \mathbb{N}$  ,  $P(x) = Q(x) + \ell x^{i_0}$  ,  $Q(x) = x^n + \dots + b x^{i_0} + \dots + b_n$  , and  $k = n - i_0$  . Then for an arbitrary  $\varepsilon > 0$ , for all

$$b \geq \left( \frac{\ell \pi \operatorname{csc}(\pi/k)}{\varepsilon k} \right)^{k/(k-1)} \tag{8}$$

the condition  $\mathcal{I} < \varepsilon$  holds for the integral (4).

*Proof.* According to the given conditions,

$$\begin{aligned} \ln \frac{P(x)}{Q(x)} &= \ln \frac{Q(x) + \ell x^{i_0}}{Q(x)} = \ln \left( 1 + \frac{\ell x^{i_0}}{Q(x)} \right) \\ &\leq \frac{\ell x^{i_0}}{Q(x)} \leq \frac{\ell x^{\ell i_0}}{x^n + b x^{i_0}} = \frac{\ell}{x^k + b} . \end{aligned}$$

Note that the first inequality is a consequence of the inequality  $\ln(1 + x) \leq x$  for  $x \geq 0$ , whereas the second of that fact that the value of a fraction increases when the nominator is decreased. Therefore,

$$\int_0^{+\infty} \ln \frac{P(x)}{Q(x)} dx \leq \ell \int_0^{+\infty} \frac{dx}{x^k + b}$$

which combined with Lemma 1 implies

$$\int_0^{+\infty} \ln \frac{P(x)}{Q(x)} dx \leq \frac{\ell}{b^{(k-1)/k}} \frac{\pi}{k} \csc \frac{\pi}{k} \leq \varepsilon$$

whenever the parameter  $b$  satisfies the condition (8). □

\* \* \* \* \*

At this point it should be noted that in the above considerations we were not “hunting” for pairs of polynomials  $P(x)$  and  $Q(x)$  with integer coefficients, for which the value of the integral  $\mathcal{I}$ , Eq. (4), assumes the smallest non-zero value. Even smaller values must be encountered if some coefficients of  $Q(x)$  are set smaller and some other greater than the respective coefficients of  $P(x)$ . Therefore, Theorem 1 may be understood as the analysis of the simplest case. Yet, its main implications are certainly valid also in the general case.

## 4 A discrete–mathematical caveat

At the first glance, from Eq. (3) and Theorem 1 it follows that there are pairs of non-equienergetic (finite) graphs whose energies differ arbitrarily little. However, Theorem 1 should be interpreted in a bit more cautious manner. In Theorem 1 it is required that the coefficient  $b$  be sufficiently large. On the other hand, in graphs with a fixed value  $n$  of vertices, the coefficients of the characteristic polynomial cannot assume arbitrarily large

values. For instance, for a bipartite graph with  $n = 2m$  vertices and maximal vertex degree  $\Delta$ , the coefficient  $b(G, k)$  in Eq. (1) is bounded by above as

$$b(G, k) \leq \Delta^{2k} \binom{m}{k}$$

with equality if and only if  $G$  consists of  $m$  isolated edges.

Anyway, because the number of graphs with a fixed value  $n$  of vertices is finite, there exists a smallest non-zero value that the energy difference does assume, say  $\varepsilon_n$ . In other words, there are no two non-equienergetic graphs of order  $n$ , such that their energies differ by less than  $\varepsilon_n$ . Consequently, for any finite  $n$ , the energy difference cannot become arbitrarily small.

Therefore, Theorem 1 should be interpreted as an indication that the energy difference may be much smaller than the earlier proposed limit  $10^{-8}$ . Since the coefficients of the characteristic polynomial rapidly increase with  $n$ , Theorem 1 also implies that very small (non-zero) energy differences are expected to be encountered at graphs with large  $n$ -values.

*Acknowledgement.* The first author was supported in part by the Serbian Ministry of Education, Science and Technological Development (grant No. 174015).

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